

Challenge to the Donor-Acceptor-Pair Photoluminescence Mechanism in Alloyed Ag-In-Zn-S Quantum Dots

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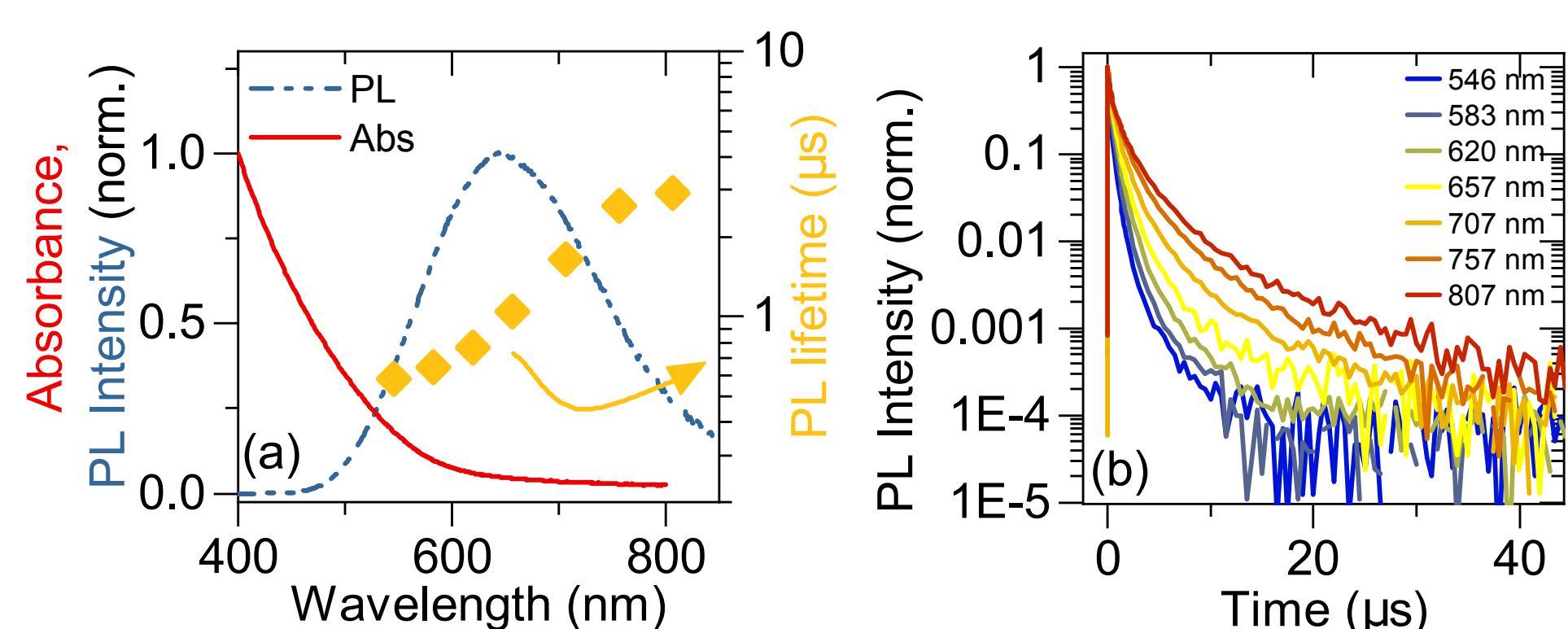
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Photoluminescence (PL)

Solution-phase PL

- Featureless absorption spectra
- Large PL linewidths (~ 0.6 eV)
- Large Stokes shifts (~ 0.5 eV)
- Long decay times (~ 500 ns)
- Wavelength-dependent PL dynamics



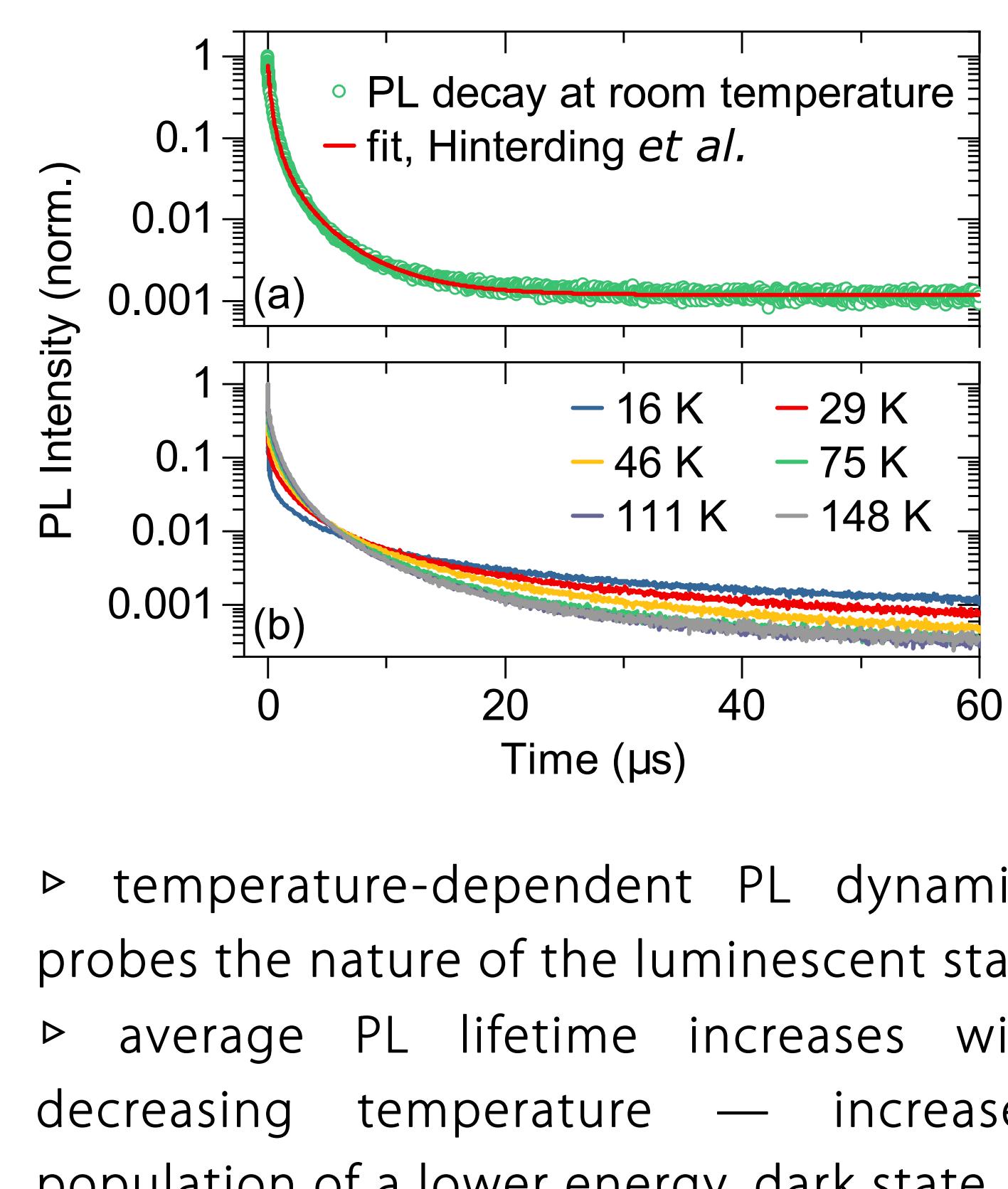
Moodley et al. *J. Mat. Chem., C* 2019, 7 11665.

Samples

Sample	Shape	Size	Structure	PL QY	Image
Ag _{1.0} In _{1.5} Zn _{1.9} S _{3.6}	Spherical	D = 3.7 nm	Orthorhombic	40%	
Ag _{1.0} In _{1.5} Zn _{4.4} S _{6.8}	Rod-like	D = 3.9 nm L = 9.0 nm	Orthorhombic	18%	

Kowalik et al. *Chem. Mater.*, 2022, 34 809.

Temperature-dependent PL dynamics



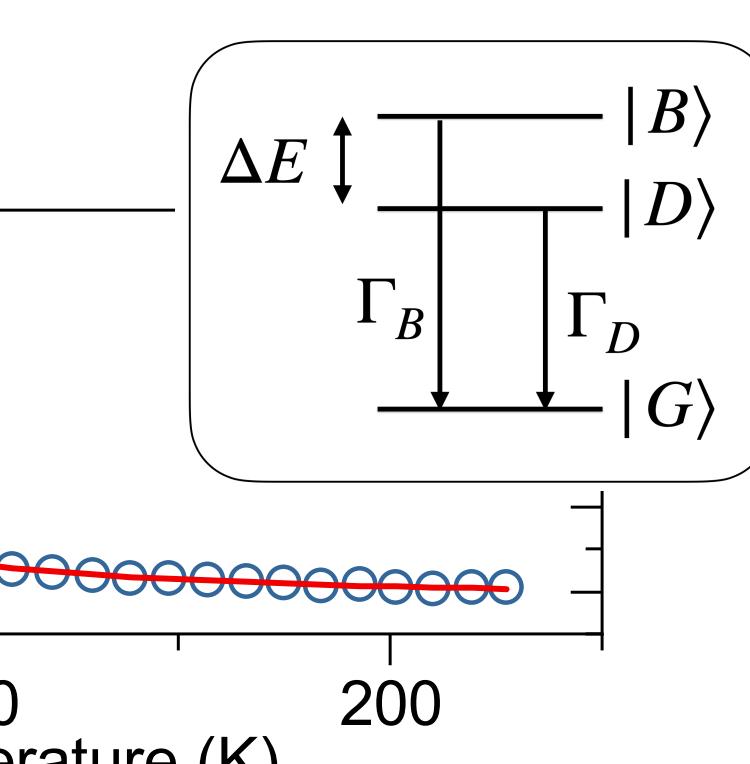
- non-exponential PL decays can be modeled by assuming electron transitions to holes at localization sites randomly distributed inside the QD volume

$$I(t) = A \int_0^{r_{\max}} \frac{\exp(-t/\tau(r_{\text{loc}}))}{\tau(r_{\text{loc}})} dr_{\text{loc}} + I_0$$

$$\tau(r_{\text{loc}}) = \tau_0 \frac{\pi^2 r_{\text{loc}}^2}{a^2 \sin^2(\pi r_{\text{loc}}/a)}$$

Hinterding et al.

Nano Lett., 2021, 21 658.



Discussion & Conclusions

- within DAP mechanism: carrier localization, strong electron-phonon coupling, and a correlation between PL energy and donor-acceptor distance explain the solution-phase PL properties

$$E_{\text{DAP}} = E_g + E_e + E_h - E_{e,\text{loc}} - E_{h,\text{loc}} + \frac{e^2}{4\pi\epsilon_0\epsilon r_{DA}}$$

- however, presence of long-lived delocalized carriers and lack of ns PL rise time, show that DAP is not the luminescent state

- our results are consistent with the PL occurring as a result of a free-to-bound transition — recombination of a delocalized electron and localized hole, similar to the case of CuInS₂ QDs.

- electron-hole exchange induces a splitting of the luminescent state and a strong temperature dependence of the PL lifetimes

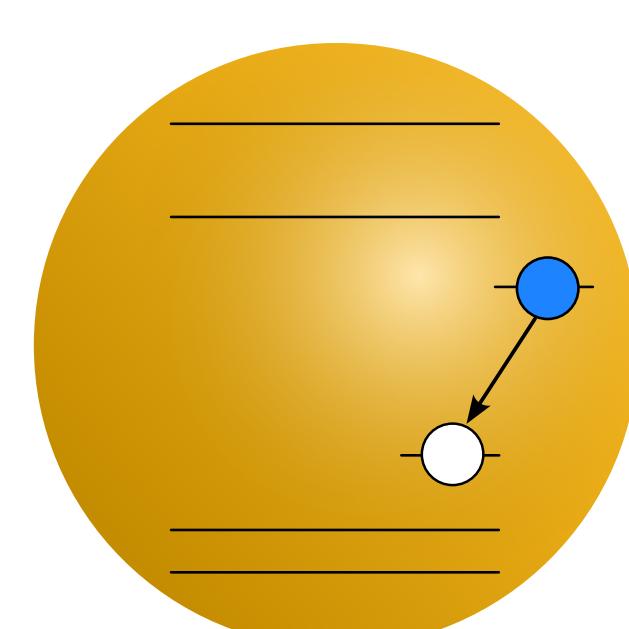
Knowles et al. *JACS*, 2015, 137 13138.
Szymura et al. *J. Phys. Chem. C*, 2023, 127, 6768.

- assuming thermal equilibrium within the fine structure and pure-spin states:

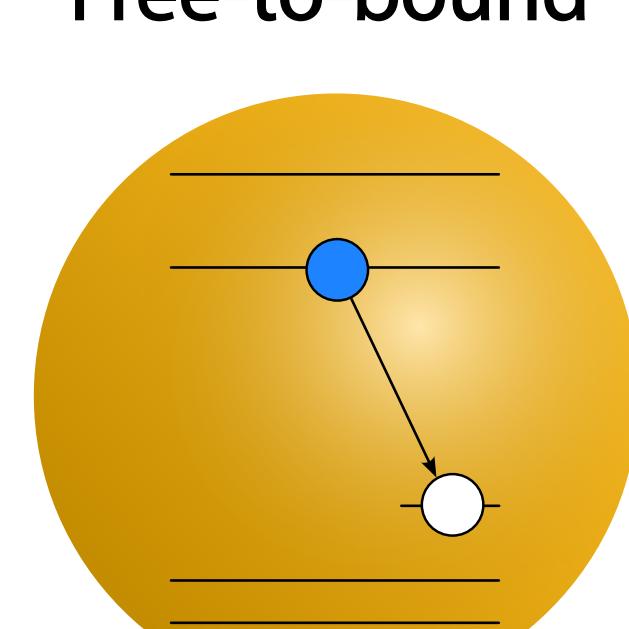
$$\tau_0(T) = \frac{1 + 3 \exp\left(\frac{\Delta E}{k_B T}\right)}{\Gamma_A + \Gamma_F 3 \exp\left(\frac{\Delta E}{k_B T}\right)}$$

- allows to retrieve the splitting at the bright state radiative lifetime for the centrally localized hole: $\Delta E = 6.3$ meV, $\tau_{0B} = 120$ ns

Donor-Acceptor Pair



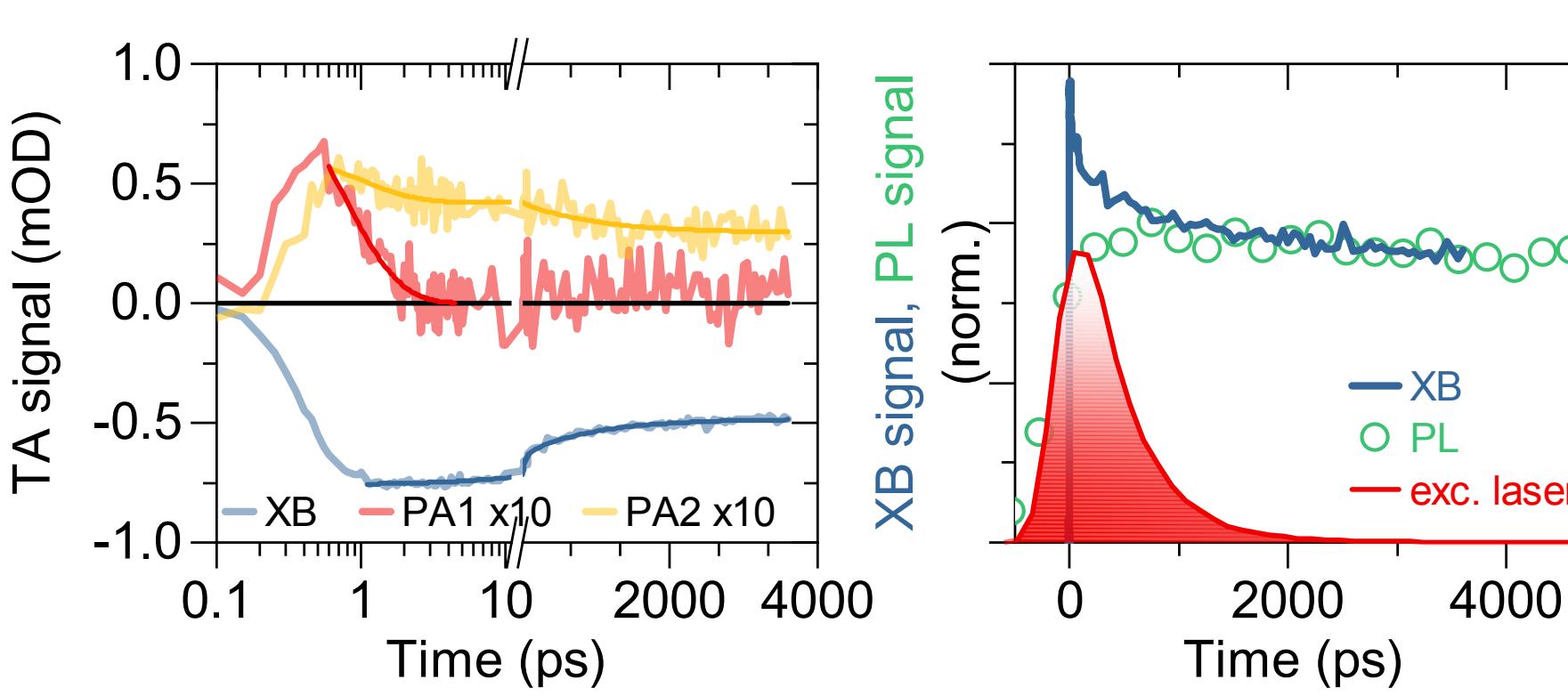
Free-to-bound



Transient absorption (TA)

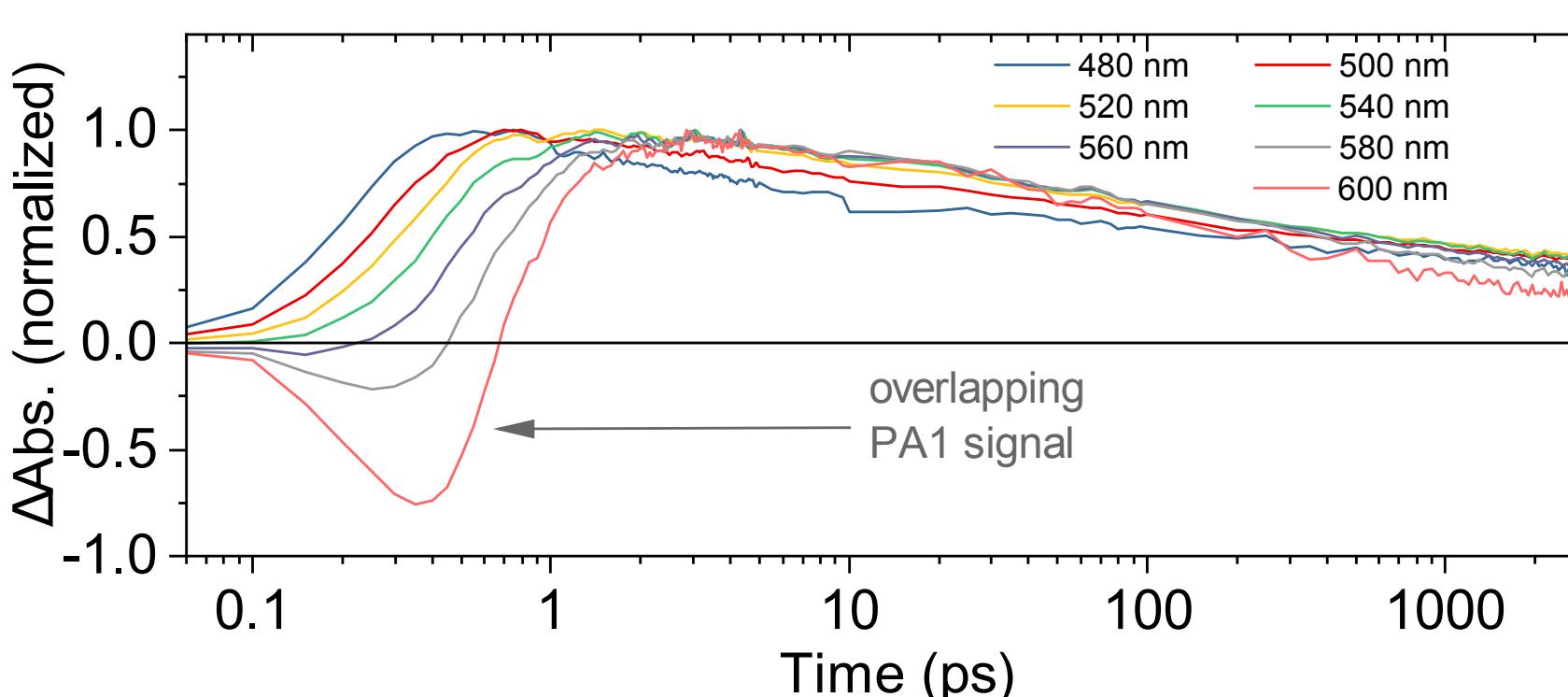
Bare Ag-In-Zn-S QDs

- XB — ground state bleach (delocalized carriers)
- PA1 — short-lived (~ 1 ps) photoinduced absorption (hot biexcitons)
- PA2 — long-lived photo-induced absorption (mid-gap localized holes)



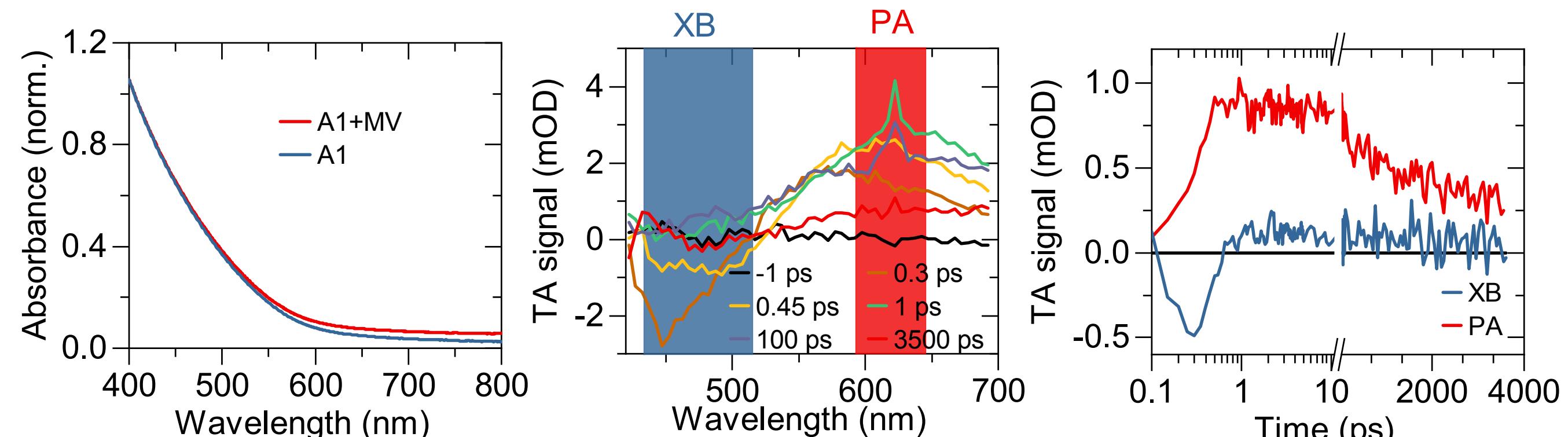
- XB, PA2 initial decay due to surface trapping
- long decay of on a timescale of tens of ns
- PL rise time on ps timescale
- conclusion: delocalized carriers participate in the PL

Carrier cooling dynamics



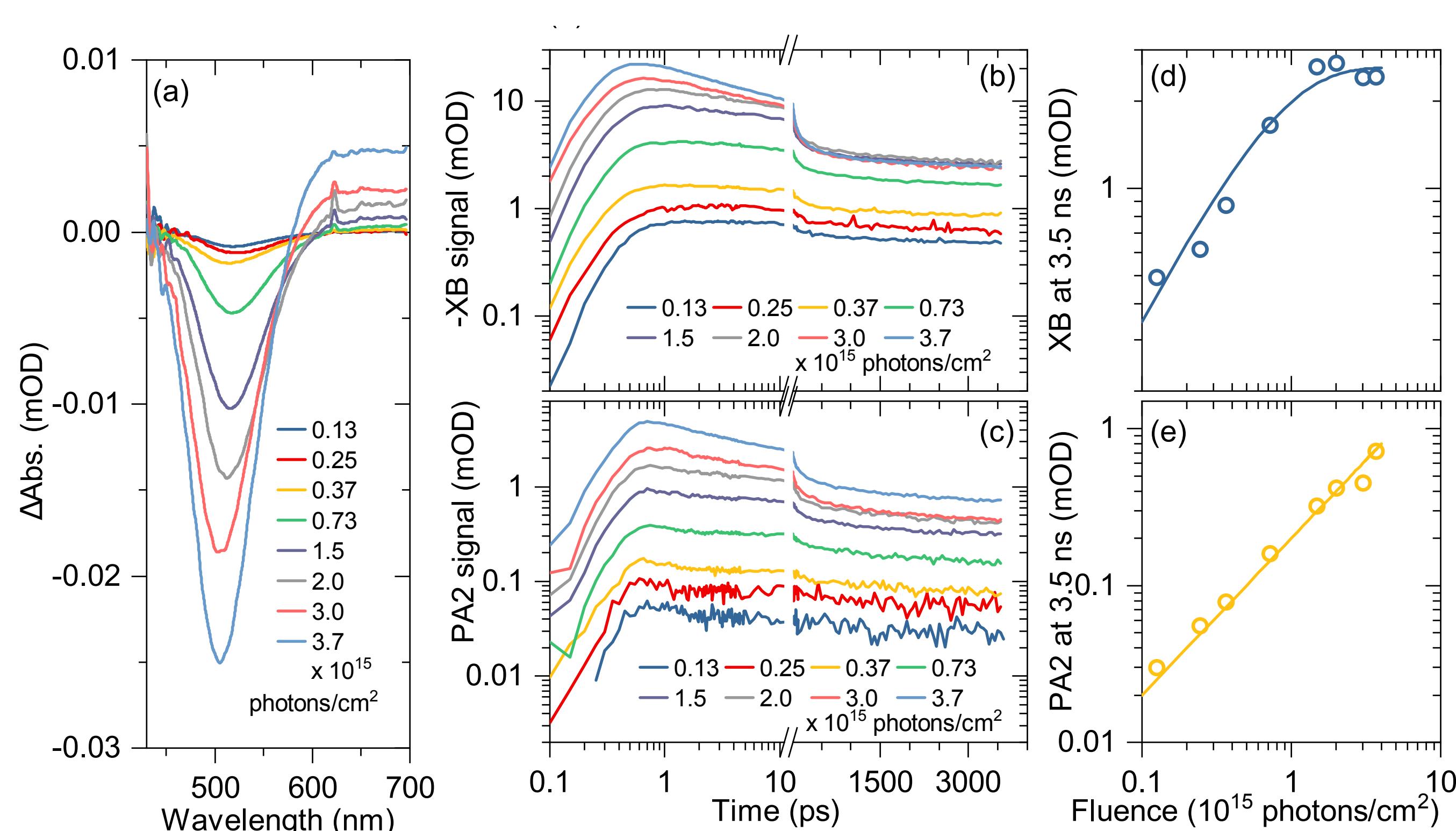
- spectral dependence of XB rise dynamics
- faster rise at shorter wavelengths (higher energies)
- carrier cooling to the ground state occurs in ~ 1 ps

QDs conjugated with electron scavenger (methyl viologen)



- conjugation with MV — no change in absorption, complete quenching of the PL
- ultrafast decay of XB signal before cooling complete — hot carrier extraction
- PA2 weakly affected, indicating its origin related to localized holes

Excitation power dependent TA



- at large delays, saturation of XB signal with excitation power — Auger recombination of multiexcitons
- XB(3.5 ns) = C(1-exp(-σJ)), where J is fluence and σ the absorption cross-section
- fitting yields $\sigma = 1.42 \times 10^{-15}$ cm², roughly 10 times smaller than for CdSe QDs
- PA2 increases linearly with power — no Auger process for localized holes